

## § Geodesics on Surfaces in $\mathbb{R}^3$

Recall: A curve  $\alpha: I \rightarrow S \subseteq \mathbb{R}^3$  is a **geodesic**

iff  $\nabla_{\alpha', \alpha'} = 0$

i.e.  $(\alpha'')^T = \nabla_{\alpha', \alpha'} = 0$ .

In local coordinates, it can be expressed a system of non-linear 2<sup>nd</sup> order ODEs:

$$\frac{d^2 u^k}{dt^2} + T^k_{ij} \frac{du^i}{dt} \frac{du^j}{dt} = 0 \quad — (*)$$

Fundamental Theorem of geodesics:

(\*) is uniquely solvable (on a short time interval)

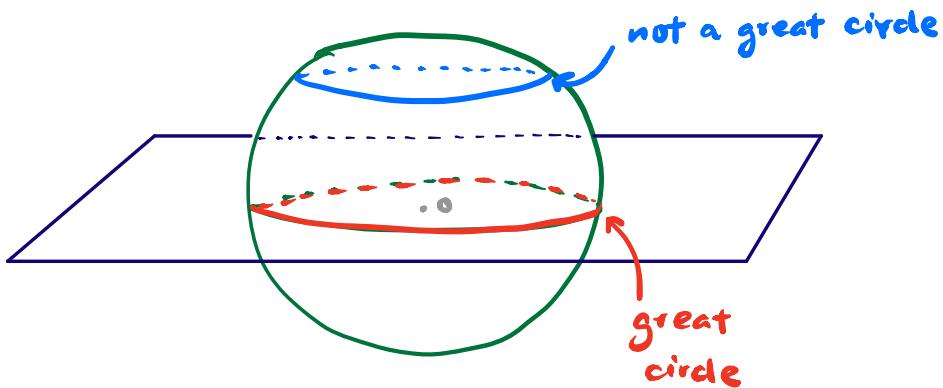
with any prescribed initial position and initial velocity.

Remark: However, explicit solutions can be very hard to compute analytically. We now look at a few simple examples that allow us to simplify the calculations by making use of the **Symmetries**!

### Example 1 : Geodesics on round spheres.

Let  $S = S^2 = \{x^2 + y^2 + z^2 = 1\}$  be the **unit sphere**.

Prop: The geodesics on  $S^2$  are given exactly by (segments) of the "**great circles**", i.e. circles obtained by intersecting  $S^2$  with a plane passing through the origin.



Proof: Suppose  $\alpha(s) : I \rightarrow S^2$  is a **geodesic p.b.a.l.**

$$\alpha \text{ lies on } S^2 \Rightarrow \|\alpha\|^2 \equiv 1 \quad \dots \dots (1)$$

$$\alpha \text{ p.b.a.l.} \implies \|\alpha'\|^2 \equiv 1 \quad \dots \dots (2)$$

$$\alpha \text{ geodesic} \implies (\alpha'')^T \equiv 0 \quad \dots \dots (3)$$

Differentiate (1),  $\langle \alpha, \alpha' \rangle \equiv 0$

Differentiate this again and use (2) :

$$\langle \alpha, \alpha'' \rangle + \underbrace{\langle \alpha', \alpha' \rangle}_{\stackrel{'''}{1}} \equiv 0$$

$$\Rightarrow \langle \alpha'', \alpha \rangle \equiv -1$$

Recall that:  $T_p S^2 \perp p$

$$\begin{aligned} \therefore \alpha'' &= (\alpha'')^T + (\alpha'')^N \\ &= \underbrace{(\alpha'')^T}_{\stackrel{\text{geodesic}}{0}} + \underbrace{\langle \alpha'', \alpha \rangle \alpha}_{\stackrel{''}{-1}} \end{aligned}$$

Hence, we arrive at the equation:

$$\alpha'' + \alpha \equiv 0$$

— (#)

Note:  $\alpha(s) = p \cos s + q \sin s$  \* solves (#)

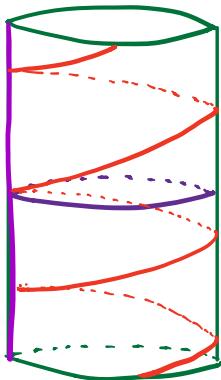
for any  $p, q \in S^2$  s.t.  $p \perp q$  (Exercise: Check this)

By choosing  $p, q$  appropriately, it solves (#) with any given initial position  $\alpha(0)$  and initial velocity  $\alpha'(0)$ .

Hence, by uniqueness, these are all the possible solutions to the geodesic equation. Finally, notice that \* parameterizes a great circle lying in the plane ( $\& S^2$ ) spanned by  $p$  and  $q$ .

## Example 2 : Geodesics on a cylinder

Let  $S = \{x^2 + y^2 = 1\}$  be a right circular cylinder (of radius 1)



Prop: The geodesics are given by segments of

- either • horizontal circle
- vertical line
- helix

We will give 2 proofs of this.

### Proof 1 : (Make use of symmetry)

As before, let  $\alpha : I \rightarrow S$  be a geodesic p.b.a.l.

Suppose  $\alpha(s) = (x(s), y(s), z(s))$ ,  $s \in I$

$$\alpha \text{ lies on } S \Rightarrow x(s)^2 + y(s)^2 \equiv 1$$

$$\xrightarrow{\text{differentiate}} xx' + yy' \equiv 0$$

$$\xrightarrow{\text{differentiate}} xx'' + yy'' + \underbrace{(x')^2 + (y')^2}_{= 1 - (z')^2} \equiv 0. \quad (\because \text{p.b.a.l.})$$

Hence, 
$$xx'' + yy'' = (z')^2 - 1 \quad \blacksquare \quad (\# \#)$$

Recall:  $T_{(x,y,z)} S \perp (x, y, 0)$

geodesic equation :  $(\alpha'')^T \equiv 0 \Rightarrow (x'', y'', z'') \parallel (x, y, 0)$

In other words,  $\exists$  function  $\lambda(s)$  s.t.

$$\left\{ \begin{array}{l} x''(s) = \lambda(s)x(s) \\ y''(s) = \lambda(s)y(s) \\ z''(s) = 0 \end{array} \right. \quad \text{--- (**)} \quad$$

Now, we solve (\*\*\*) with initial conditions:

$$\alpha(0) = (x(0), y(0), z(0)) = (1, 0, 0)$$

$$\alpha'(0) = (x'(0), y'(0), z'(0)) = (0, a, b) \perp \alpha(0)$$

$$\text{where } a^2 + b^2 = 1 \quad \text{p.b.a.l.}$$

Solving for  $z$ , we have  $z(s) = bs \quad (\Rightarrow z' \equiv b)$

Put everything back into (##)

$$\lambda(s) = x(s)x''(s) + y(s)y''(s) = b^2 - 1 = -a^2.$$

( $\because x^2 + y^2 \leq 1$ )

constant!

Solve  $x'' = -a^2 x$  with  $x(0) = 1, x'(0) = 0$

$$x(s) = \cos as$$

Solve  $y'' = -a^2 y$  with  $y(0) = 0, y'(0) = a$

$$y(s) = \sin as$$

In summary, we have

$$\alpha(s) = (\cos as, \sin as, bs) \quad s \in I$$

where  $a, b \in \mathbb{R}$  are constants s.t.  $a^2 + b^2 = 1$ .

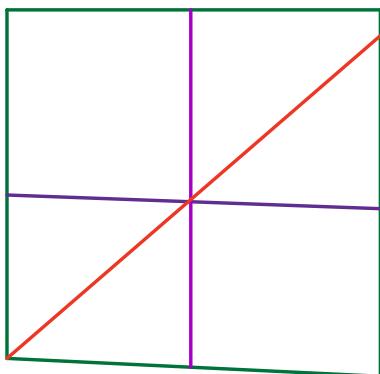
Case 1:  $b=0, a=1 \Rightarrow$  horizontal circle

Case 2:  $b=1, a=0 \Rightarrow$  vertical line

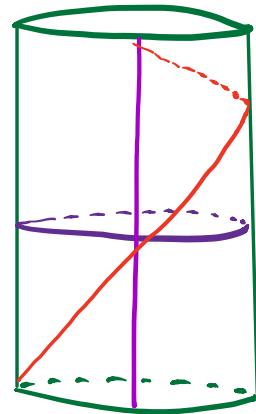
Case 3:  $b \neq 0, a \neq 0 \Rightarrow$  helix

————— □

Proof 2: Geodesics are intrinsic concepts, thus is preserved by (local) isometries.



"wrap around"  
local isometry



————— □

## § First Variation Formula for arc length

Let  $\alpha : [0, L] \rightarrow S$  be a curve on  $S$  p.b.a.l.

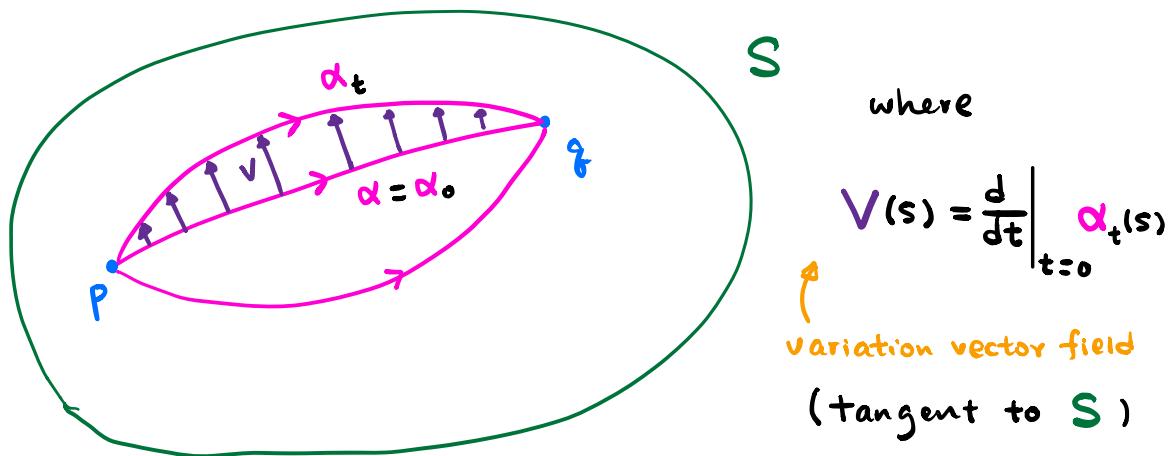
with end points  $p = \alpha(0), q = \alpha(L)$ .

Consider a "variation" of  $\alpha$  with fixed endpoints:

$$\left\{ \alpha_t(s) \right\}_{t \in (-\varepsilon, \varepsilon)} \text{ s.t. } \alpha_0 = \alpha$$

$\alpha_t : [0, L] \rightarrow S$  and  $\alpha_t(0) = \alpha_0(0)$  for all  $t \in (-\varepsilon, \varepsilon)$

$$\alpha_t(L) = \alpha_0(L)$$



First Variation Formula :

$$\frac{d}{dt} \Big|_{t=0} \text{Length}(\alpha_t) = - \int_0^L \langle V(s), \alpha''(s) \rangle ds$$

Remark: If  $\alpha$  is a critical point for ALL variations,

then

$$0 = \left. \frac{d}{dt} \right|_{t=0} \text{Length}(\alpha_t) = - \int_0^L \langle V(s), \alpha''(s) \rangle ds$$

for ALL variation field  $V$   
(which is tangential)

$$\Rightarrow (\alpha'')^T \equiv 0$$

Hence, geodesics are exactly the **critical points** to the arc length functional!

Proof of 1<sup>st</sup> Variation Formula:

By definition of arc length, (Note:  $\alpha_t(s)$  is NOT p.b.a.l.  
except for  $t=0$ )

$$\text{Length}(\alpha_t) = \int_0^L \|\alpha'_t(s)\| ds$$

Differentiate w.r.t.  $t$ .

$$\begin{aligned} \frac{d}{dt} \text{Length}(\alpha_t) &= \int_0^L \frac{1}{\|\alpha'_t(s)\|} \left\langle \frac{\partial}{\partial t} \left( \frac{\partial \alpha}{\partial s} \right), \frac{\partial \alpha}{\partial s} \right\rangle ds \\ &= \int_0^L \frac{1}{\|\alpha'_t(s)\|} \left\langle \frac{\partial}{\partial s} \left( \frac{\partial \alpha}{\partial t} \right), \frac{\partial \alpha}{\partial s} \right\rangle ds \end{aligned}$$

*switch*

Evaluate at  $t=0$ ,  $\|\alpha'_0(s)\| \equiv 1$ ,  $\left. \frac{\partial \alpha}{\partial t} \right|_{t=0} = V$

$$\begin{aligned}
 \left. \frac{d}{dt} \right|_{t=0} \text{Length}(\alpha_t) &= \int_0^L \left\langle \frac{\partial}{\partial s} V, \frac{\partial}{\partial s} \alpha \right\rangle ds \\
 &= - \int_0^L \left\langle V, \alpha'' \right\rangle ds \\
 &\quad + \left. \left\langle V, \frac{\partial}{\partial s} \alpha \right\rangle \right|_{s=0}^{s=L} = 0
 \end{aligned}$$

(fixed  
end points)